

## Pion Form Factor With Twisted Mass QCD

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## Motivation for Twisted Mass LQCD:

- Lattice QCD with Wilson fermions:

$$S_F^{LW} = a^4 \sum_x \bar{\psi}_p(x) \left[ \sum_{\mu} \frac{1}{2} \gamma_{\mu} (\nabla_{\mu}^* + \nabla_{\mu}) + m_0 - a \frac{r}{2} \sum_{\mu} \nabla_{\mu}^* \nabla_{\mu} \right] \psi_p(x)$$

- $\mathcal{O}(a)$  errors
- Exceptional configurations in quenched calculations.
- Improvement with *clover* term and improved operators ...  
**exceptional configurations** still exist.

## Twisted Mass QCD and Exceptional Configurations:

- *R. Frezzotti, P. A. Grassi, S. Sint and P. Weisz, JHEP **0108**, 058 (2001).*
- *R. Frezzotti, S. Sint and P. Weisz, JHEP **0107**, 048 (2001).*

- In the continuum, the twisted mass action for a *degenerate quark doublet* is given by:

$$S_F[\psi, \bar{\psi}] = \int d^4x \bar{\psi} (D_\mu \gamma_\mu + m_q + i\mu_q \gamma_5 \tau^3) \psi$$

- An axial transformation,

$$\psi' = e^{i\omega \gamma_5 \tau^3 / 2} \psi, \quad \bar{\psi}' = \bar{\psi} e^{i\omega \gamma_5 \tau^3 / 2}$$

leaves the form of the action invariant and rotates the masses as:

$$\begin{pmatrix} m'_q \\ \mu'_q \end{pmatrix} = \begin{bmatrix} \cos(\omega) & \sin(\omega) \\ -\sin(\omega) & \cos(\omega) \end{bmatrix} \begin{pmatrix} m_q \\ \mu_q \end{pmatrix}$$

- Quark mass:

$$M_q^2 = m_q^2 + \mu_q^2$$

- $\tan(\omega) = \frac{\mu_q}{m_q} \rightarrow \mu'_q = 0$  which is QCD.

- Composite field operators:

$$A'^a_\mu = \begin{cases} \cos(\omega) A^a_\mu + \epsilon^{3ab} \sin(\omega) V^b_\mu & (a = 1, 2), \\ A^3_\mu & (a = 3) \end{cases}$$

- Positive definite determinant of the Dirac operator  
 $\rightarrow$  **No exceptional configurations**.

## Twisted Mass and $\mathcal{O}(a)$ improvement:

- *R. Frezzotti and G. C. Rossi, Nucl. Phys. Proc. Suppl. **129**, 880 (2004)*
- *R. Frezzotti and G. C. Rossi, arXiv:hep-lat/0306014.*

- Symmetries:

$$\begin{aligned}\mathcal{R}_5^{SP} &= \mathcal{R}_5 \times (r \rightarrow -r) \times (m_q \rightarrow -m_q) \\ \mathcal{R}_5 &\times \mathcal{D}_d \\ \mathcal{R}_5 &: \psi \rightarrow \gamma_5 \psi, \quad \bar{\psi} \rightarrow -\bar{\psi} \gamma_5 \\ \mathcal{D}_d &: U_\mu(x) \rightarrow U_\mu^\dagger(-x - a\hat{\mu}) \\ &\quad \psi(x) \rightarrow e^{3i\pi/2} \psi(-x)\end{aligned}$$

- Wilson averaging gives  $\mathcal{O}(a)$  improved correlators:

$$\frac{1}{2}[\langle O \rangle_r^\omega + \langle O \rangle_{-r}^\omega] = \zeta_O(\omega, r) \langle O \rangle_{cont.} + \mathcal{O}(a^2)$$

- Wilson averaging is done with *tmQCD*.

- Simplification at  $\omega = \pm \frac{\pi}{2}$

- Matrix elements are either automatically improved or improved by averaging correlator with momenta  $\vec{k}$  and  $-\vec{k}$ .

$$\begin{aligned} \left\langle n, \vec{k} | O | n', \vec{k}' \right\rangle_{r, m_q}^{\frac{\pi}{2}} + \eta_{nn'O} \left\langle n, -\vec{k} | O | n', -\vec{k}' \right\rangle_{r, m_q}^{\frac{\pi}{2}} \\ = 2\zeta_O(r) \left\langle n, \vec{k} | O | n', \vec{k}' \right\rangle_{cont.}^{\frac{\pi}{2}} + \mathcal{O}(a^2) \end{aligned}$$

$$\eta_{nn'O} = \pm 1$$

$$E_n(\vec{k}, r, m_q) + E_n(-\vec{k}, r, m_q) = 2E_n^{cont.}(\vec{k}, m_q) + \mathcal{O}(a^2)$$

## Details of Our Lattice Calculations:

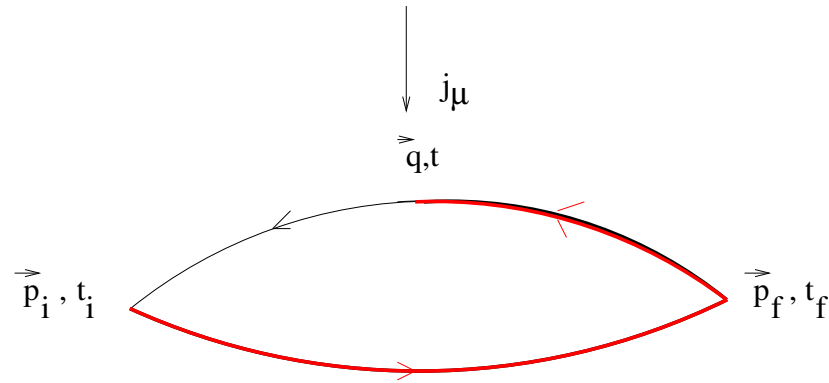
- Lattice size:  $16^3 \times 48$
- 100 gauge field configurations, quenched.
- Matrix inversion by GMRES-DR
  - *R. B. Morgan SIAM J. Sci. Comput.* **24**, 20 (2002).
  - *R. B. Morgan and W. Wilcox, Nucl. Phys. Proc. Suppl.* **106**, 1067 (2002).
- $\beta = 6.0$  ,  $\kappa = \kappa_c$ ,  $\mu_q = 0.015, 0.030$ .
  - *K. Jansen, A. Shindler, C. Urbach and I. Wetzorke, Phys. Lett. B* **586**, 432 (2004).



## The Pion Form Factor

- The pion form factor is of theoretical and experimental interest.
- No disconnected lattice diagrams.
- Transition from non-perturbative QCD to perturbative QCD occurs at smaller  $Q^2$  than for baryons.
- Mainly described by Vector Meson Dominance.
- Experimental measurement at high  $Q^2$  is underway at **JLab**.

- Various attempts to compute the pion form factor on the lattice:
  - Wilson (quenched)
    - \* *F. Bonnet, R. Edwards, G. Fleming, R. Lewis and D. Richards [LHP Collaboration], Nucl. Phys. Proc. Suppl. **128**, 59 (2004), Nucl. Phys. Proc. Suppl. **129**, 206 (2004).*
  - Clover (quenched)
    - \* *J. van der Heide, M. Lutterot, J. H. Koch and E. Laermann, Phys. Lett. B **566**, 131 (2003).*
    - \* *J. van der Heide, J. H. Koch and E. Laermann, Phys. Rev. D **69**, 094511 (2004).*
  - Domain Wall (quenched)
    - \* *Y. Nemoto [RBC Collaboration], Nucl. Phys. Proc. Suppl. **129**, 299 (2004).*
  - Domain Wall (dynamical staggered configurations)
    - \* *G. Fleming [LHP Collaboration], to be presented in this conference.*



- It corresponds to the process (for example):

$$\pi(\vec{p}_i) \rightarrow \pi(\vec{p}_f) \gamma^*(\vec{q})$$

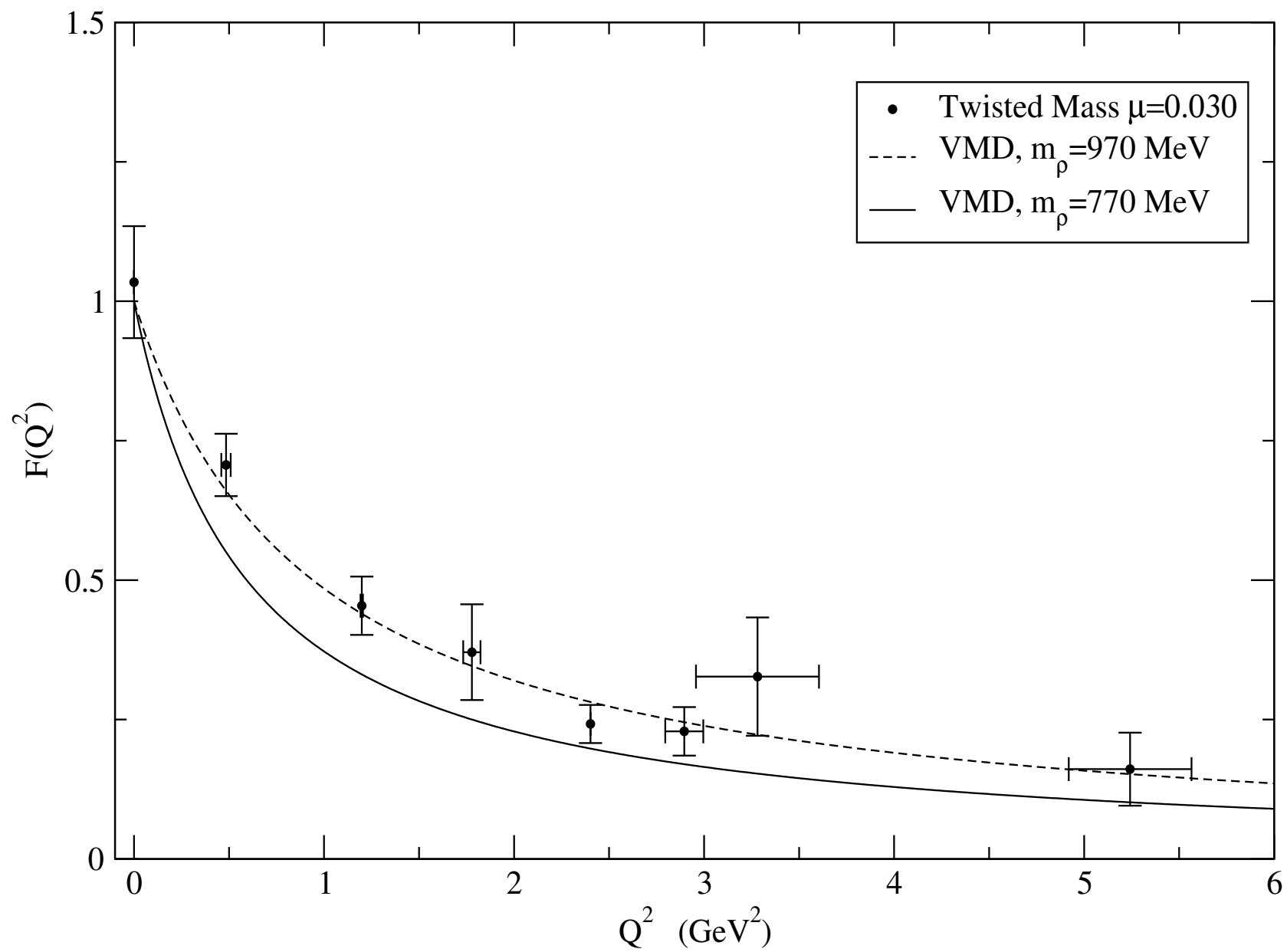
- The pion form factor  $F(Q^2)$  is defined by:

$$\begin{aligned} \langle \pi(p_f) | j_\mu(0) | \pi(p_i) \rangle_{cont.} &= F(Q^2) (p_i + p_f)_\mu \\ Q &= p_f - p_i \end{aligned}$$

- $j_\mu$  is a conserved vector current

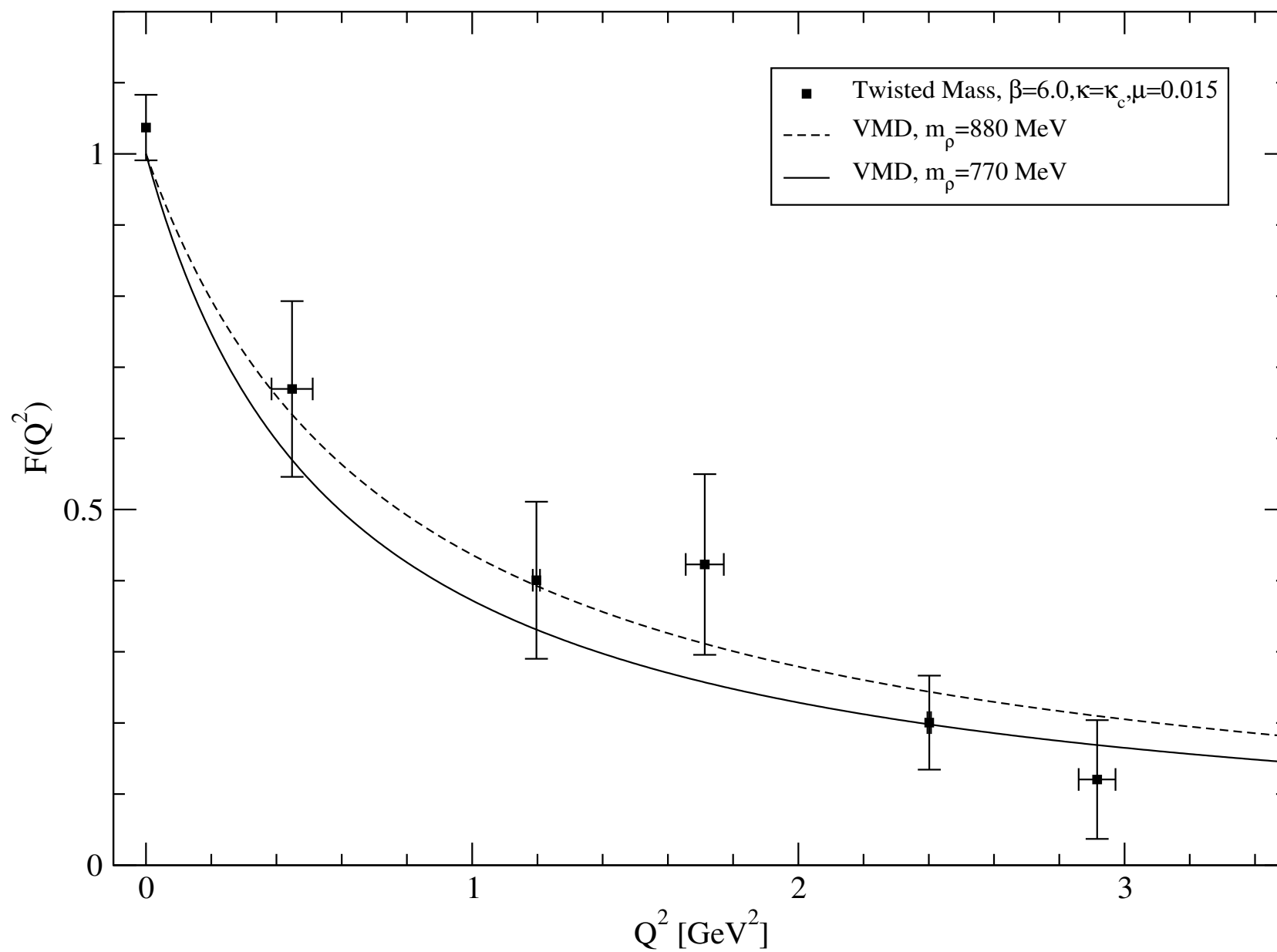
# Pion form factor at $\mu=0.030$ , $\omega=\pi/2$

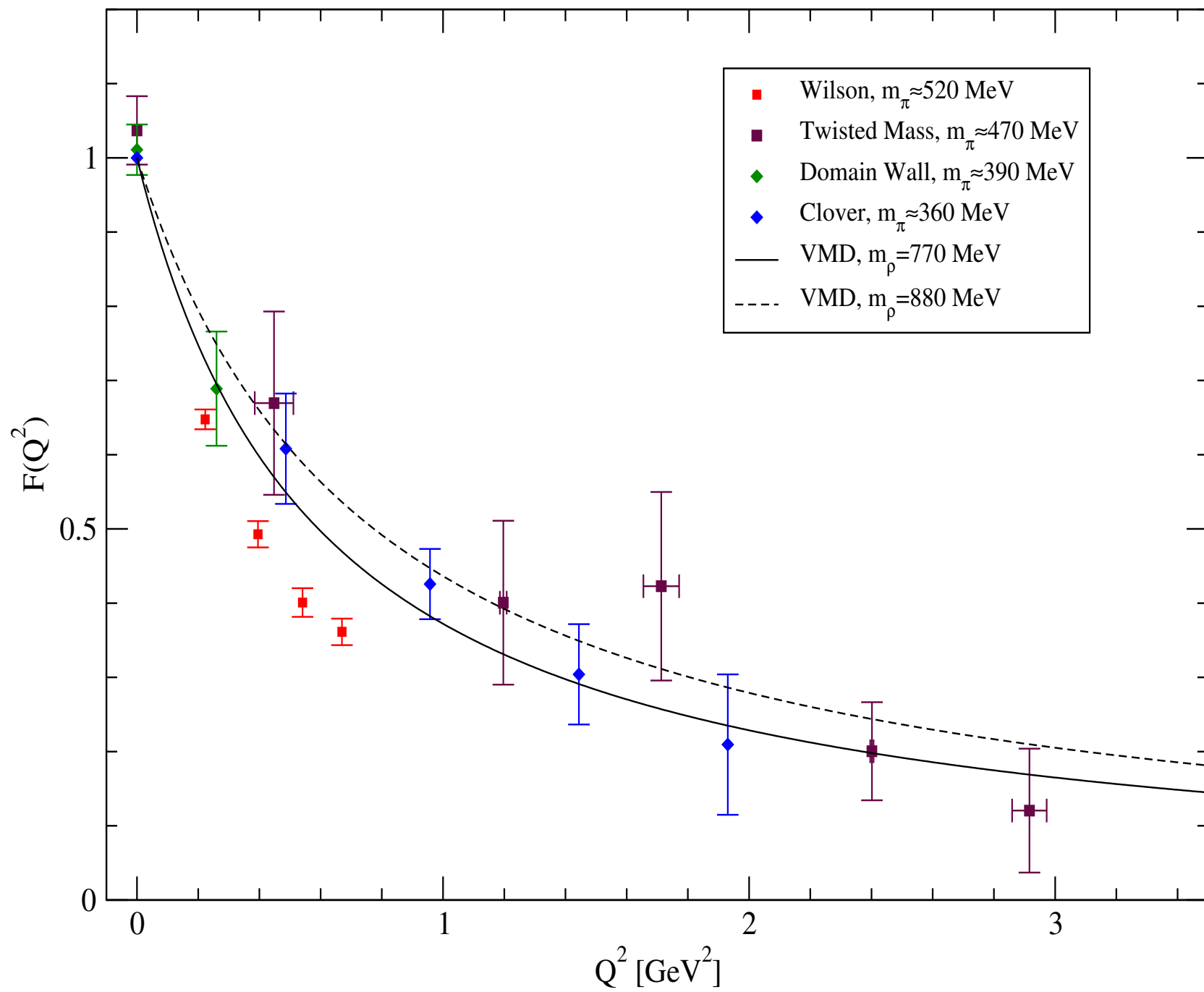
$m_\pi=660$  MeV



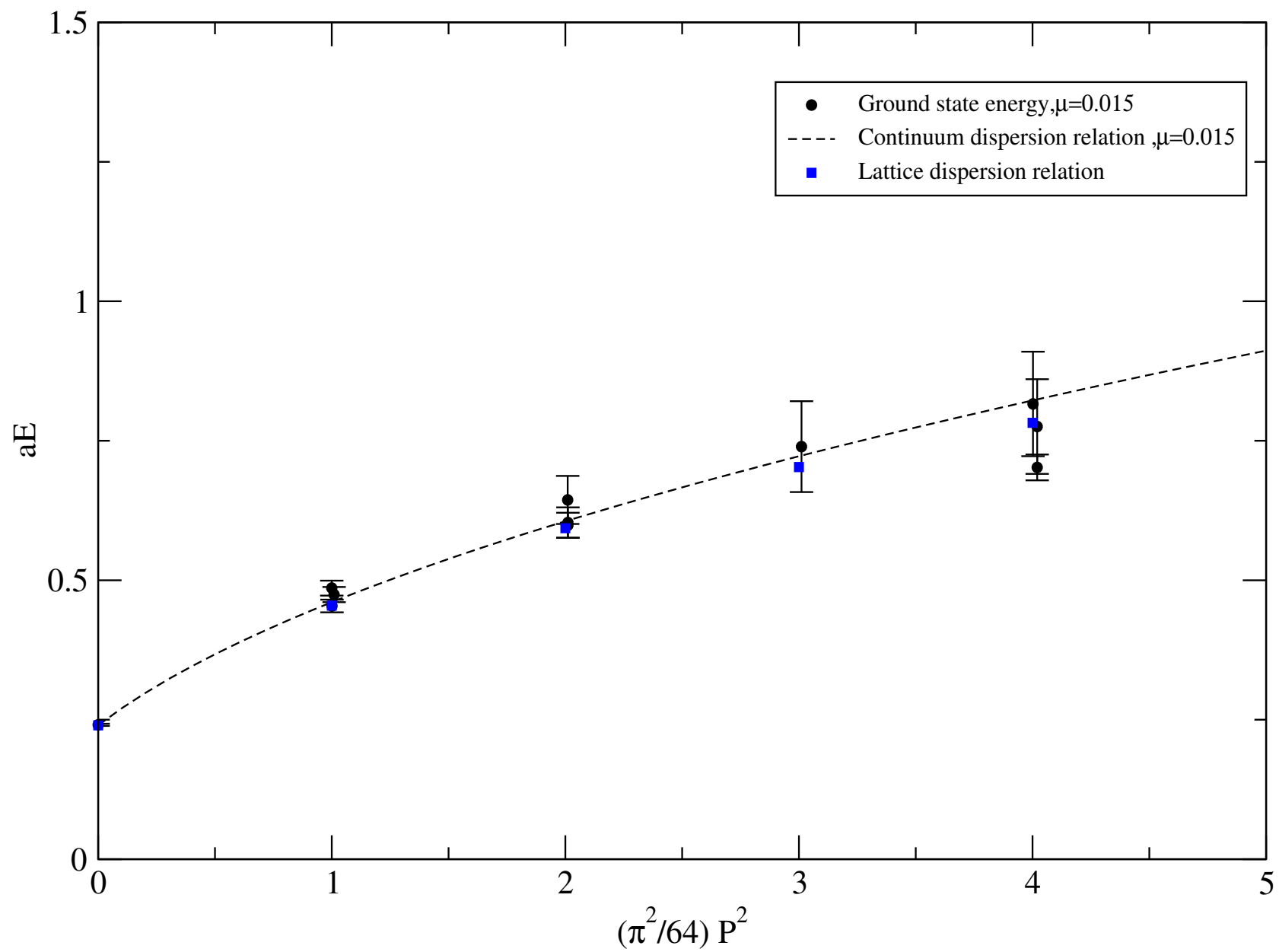
# Pion Form Factor at $\mu=0.015$ , $\omega=\pi/2$

$m_\pi=470$  MeV

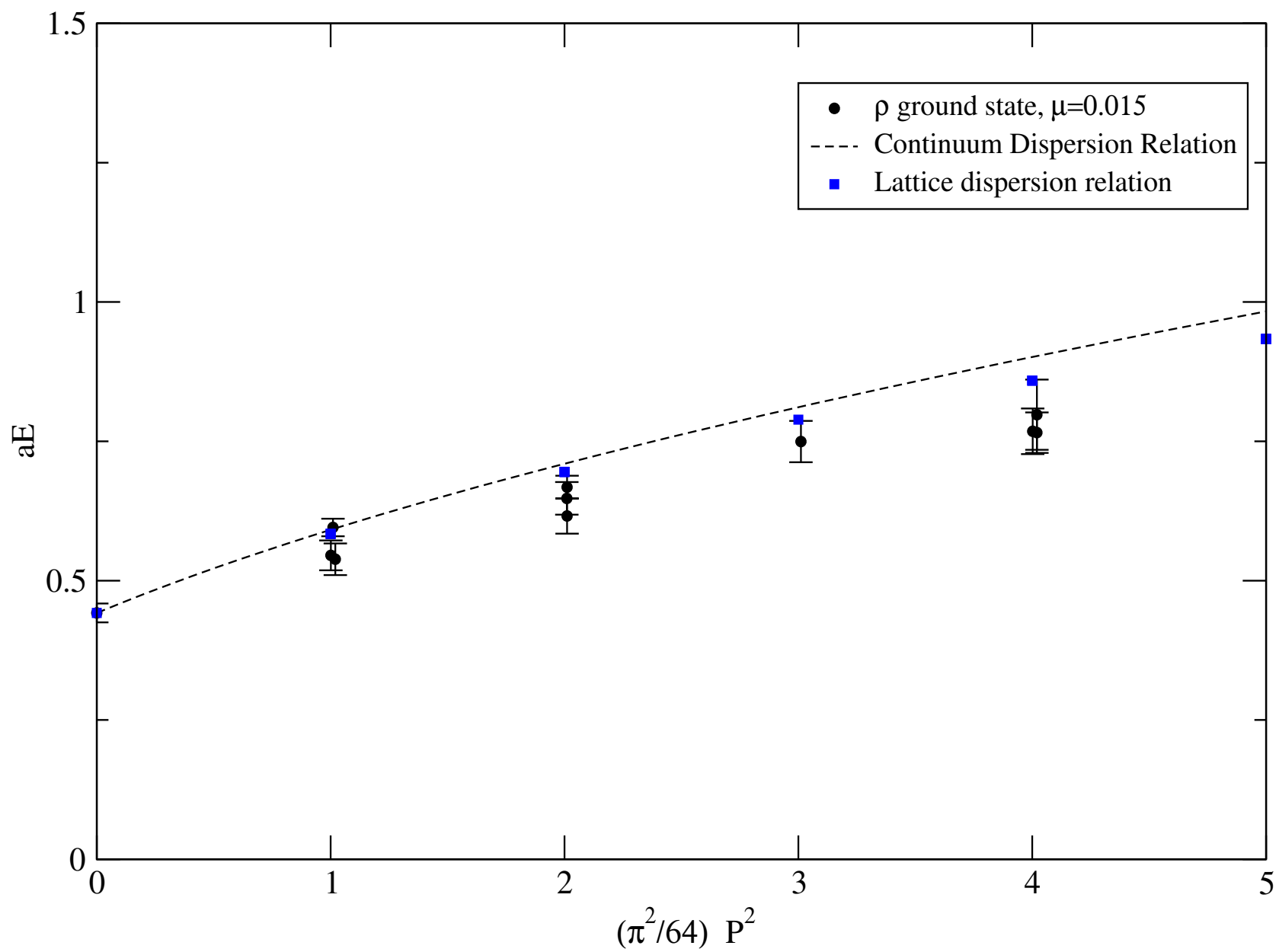




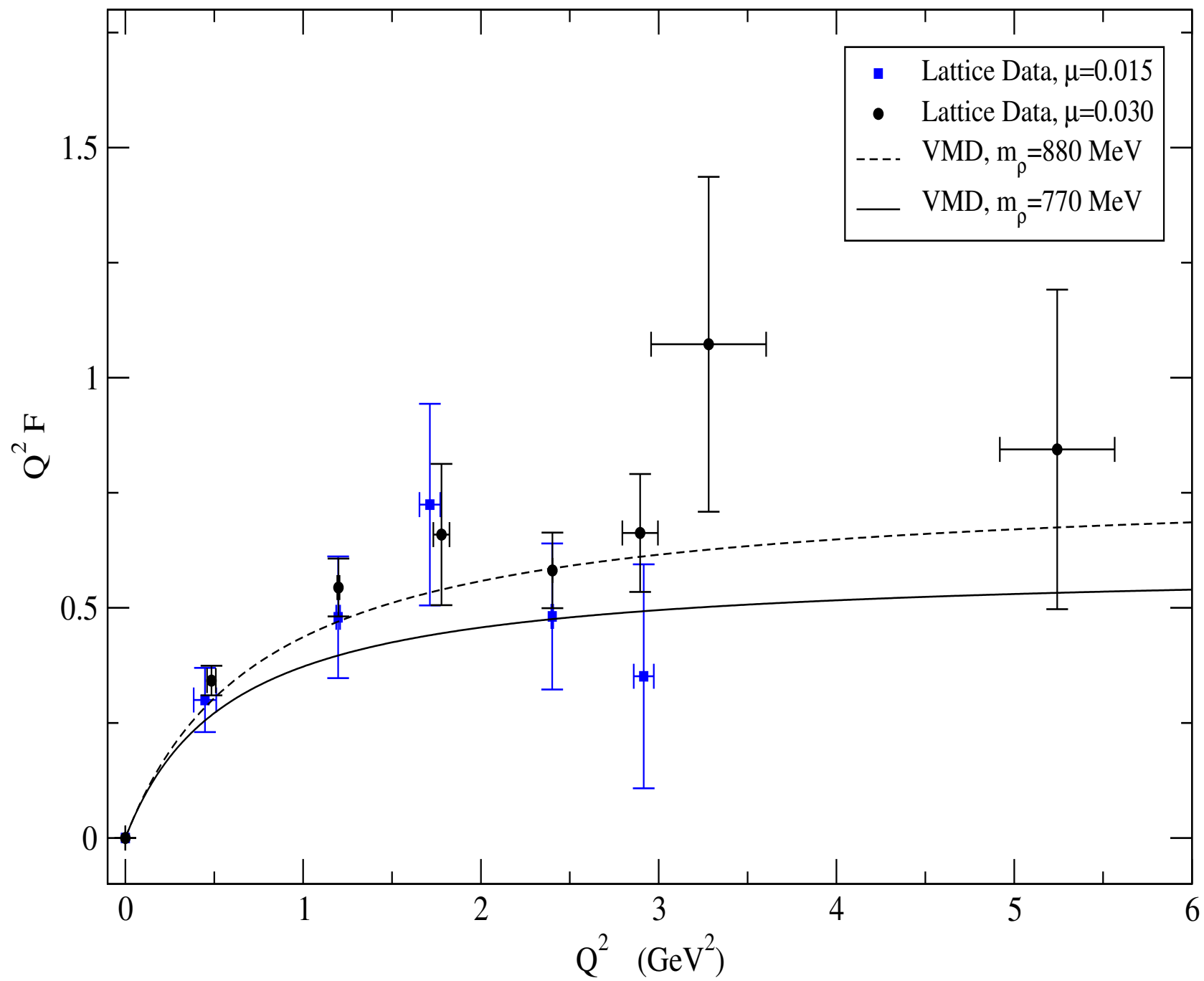
Pion Dispersion Relation at  $\mu=0.015$

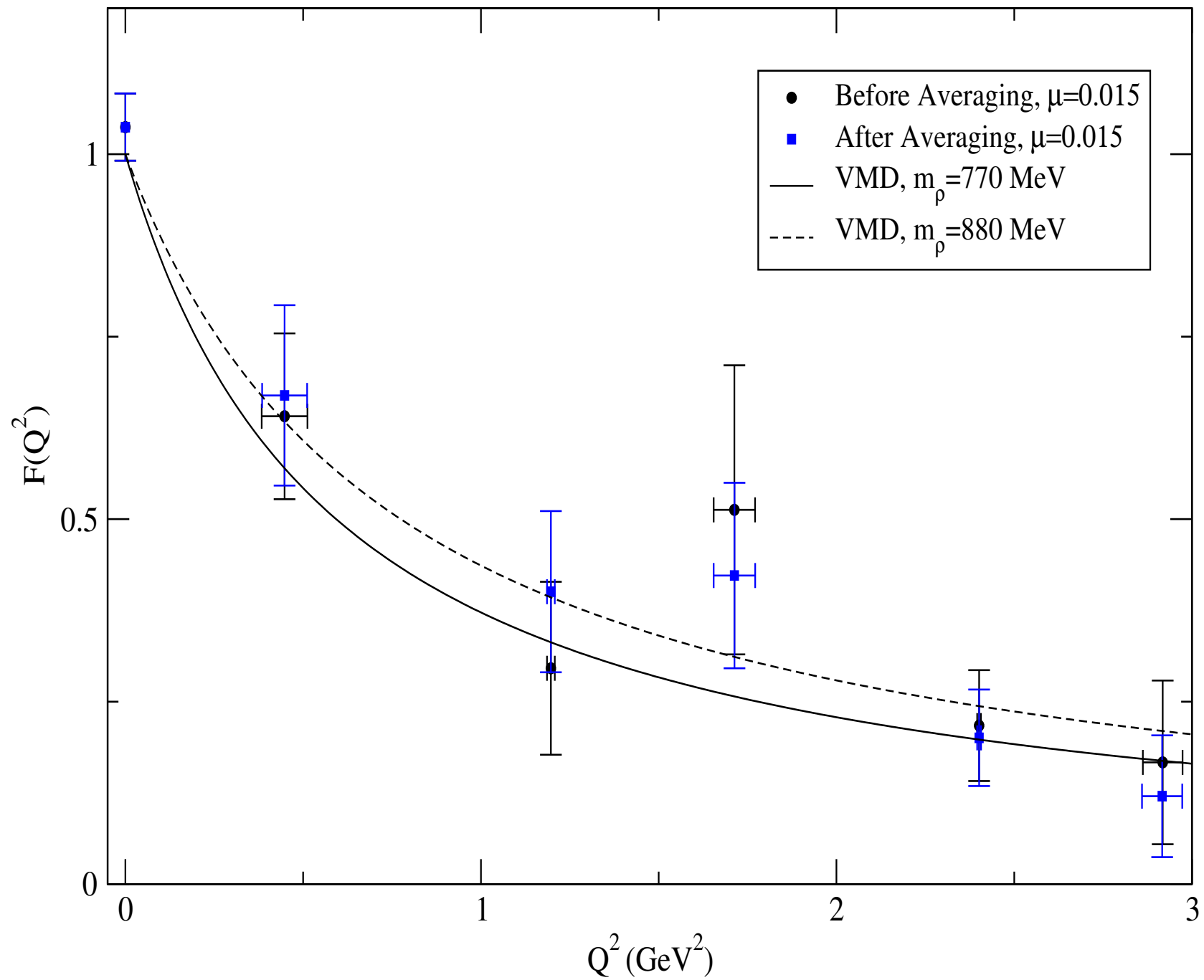


$\rho$  Dispersion Relation at,  $\mu=0.015$









## Conclusions:

- `tmQCD` solves the problem of exceptional configurations.
- $\mathcal{O}(a)$  improvement could be obtained with a simple averaging procedure at  $\omega = \pm\frac{\pi}{2}$  without the need for improvement terms.
- At  $\omega = \pm\frac{\pi}{2}$  the results are improved as compared to Wilson even before averaging.
- `GMRES-DR` is an efficient matrix inverter for *tmQCD* calculations.
- The pion form factor was calculated to a high  $Q^2$  value.